

Directions in Fundamental Limitations of Quantum Technologies

FLQT Dublin

Themes of This Talk

**Fundamental Limitations
can arise from**

Part I

**Device
Noise**

Part II

**Finite
Samples**

Part III

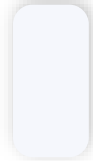
**Computational
Restrictions**

Part I

How bad can
noise really be?

Noisy Quantum Devices

Implementations of quantum operations are imperfect

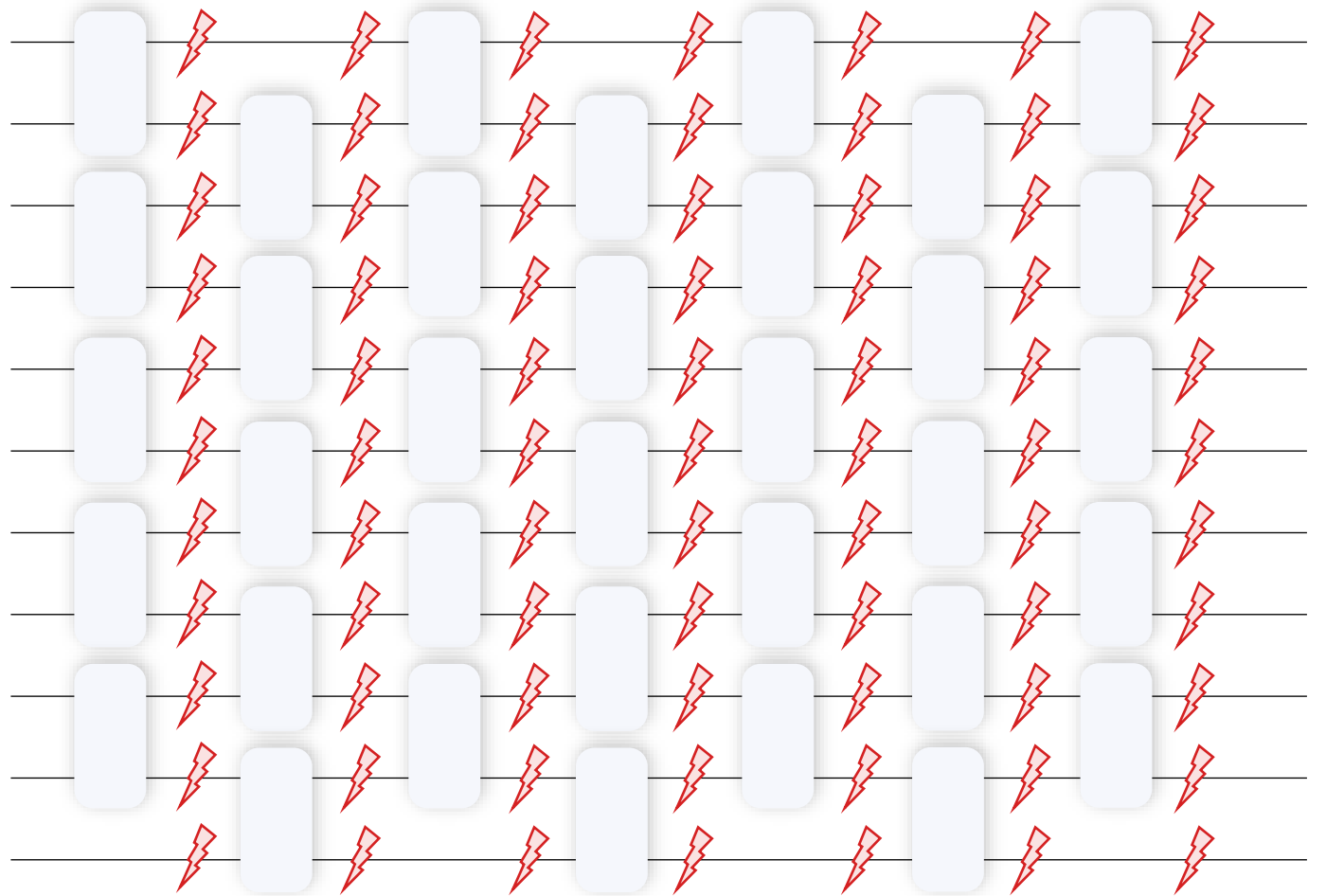


We can model this using noise channels



Depolarizing Noise

$$\mathcal{D}_\lambda[\rho] = \lambda\rho + (1 - \lambda)\omega$$



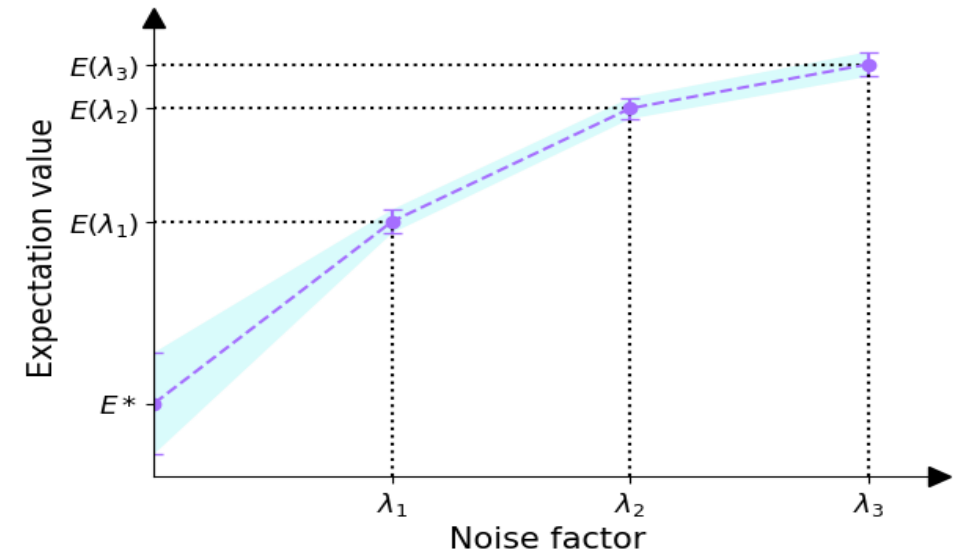
Error Mitigation

KEY IDEA

Can we gather more data and use it to **classically** „undo“ the noise?

ZERO NOISE EXTRAPOLATION¹

- › Increase noise level by adding superfluous operations
- › Compute expectation values at different noise levels
- › Fit the curve and extrapolate to zero noise



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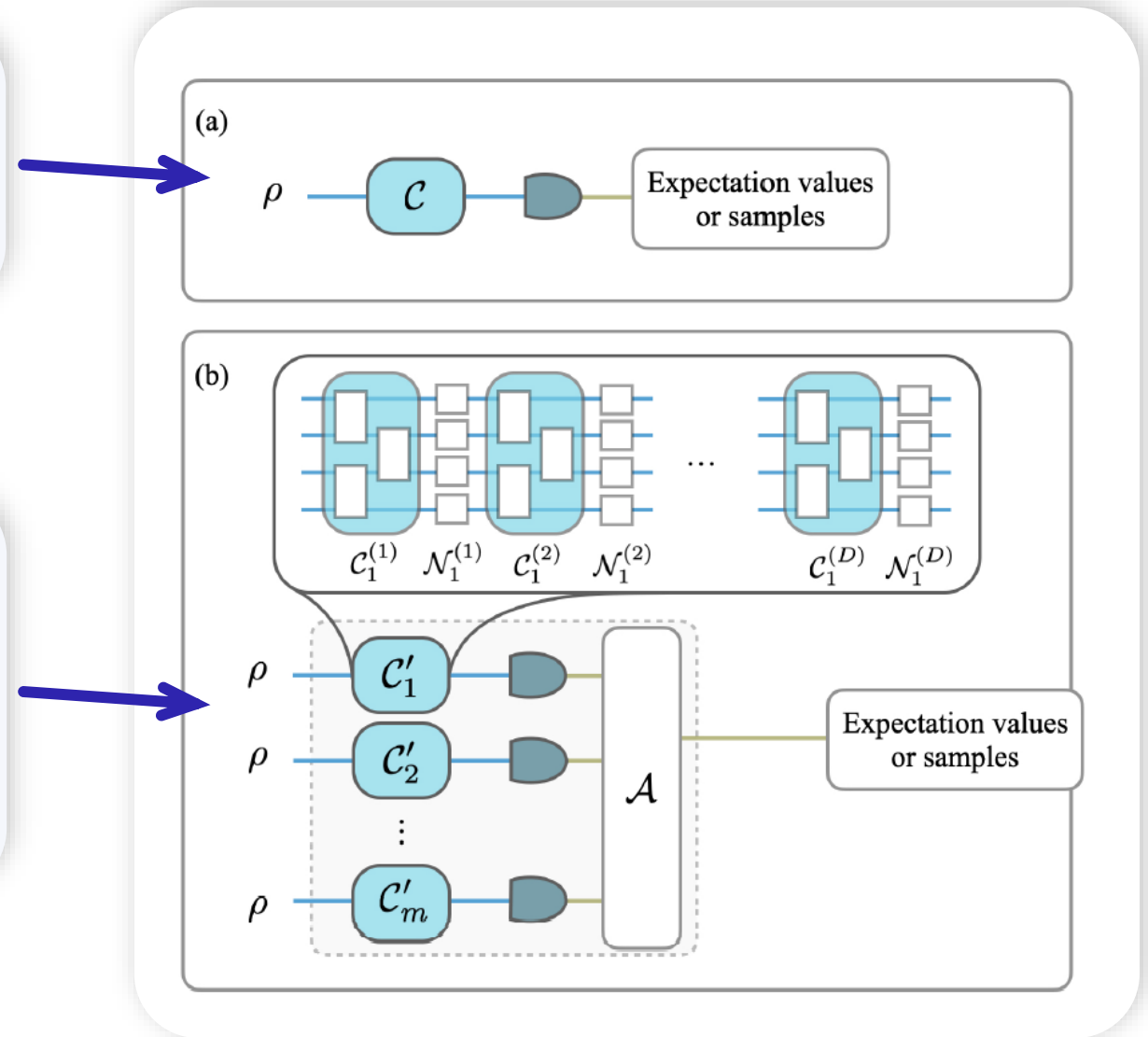
Error Mitigation

IDEAL CASE

We run a target circuit on an input state and obtain the desired output

ERROR MITIGATION ABSTRACTION¹

We run a number of noisy circuits on the input state and post-process their outputs into a prediction



¹cf. Takagi et al., npj QI 8(1), 114 (2022); Takagi et al., PRL 131(21), 210602 (2023); Tsubouchi et al., PRL 131(21), 210601 (2023).

Limitations on Error Mitigation

LOCAL DEPOLARIZING NOISE

Pushes system state towards the maximally mixed state exponentially fast in depth¹

We need exponential-in-depth resources for quantum error mitigation^{2,3}

$$N \geq \exp(O(D))$$

WE SHOW THAT

For many circuits, this effect is drastically amplified, even in log-log depth

We need exponential-in-depth and in-qubit-count resources for quantum error mitigation

$$N \geq \exp(O(nD))$$

¹Müller-Hermes et al., JMP 57(2) (2016). ²Takagi et al., PRL, 131(21), 210602 (2023) ³Tsubouchi et al., PRL 131(21), 210601 (2023).

Generic Bounds for Noisy Applications

LOCAL DEPOLARIZING NOISE

Pushes system state towards the maximally mixed state exponentially fast in depth

GENERICITY

This holds independent of the actual unitaries in the circuit

TIGHTNESS

There exist quantum circuits that saturate this scaling

Different circuits should yield different outputs (for example expectation values)

If the states for different circuits are close to the maximally mixed state they are mutually close

If outputs are far but states are close, we have to use many copies!

Beyond Generic Bounds

DECAY IN WIDTH

does not happen generically,
consider circuits that only contain
single-qubit gates

WE PROVE

That these states converge to
the maximally mixed state
much faster

WE NEED CIRCUITS

That create lots of entanglement
in short depth. We construct them
using unitary 2-designs¹

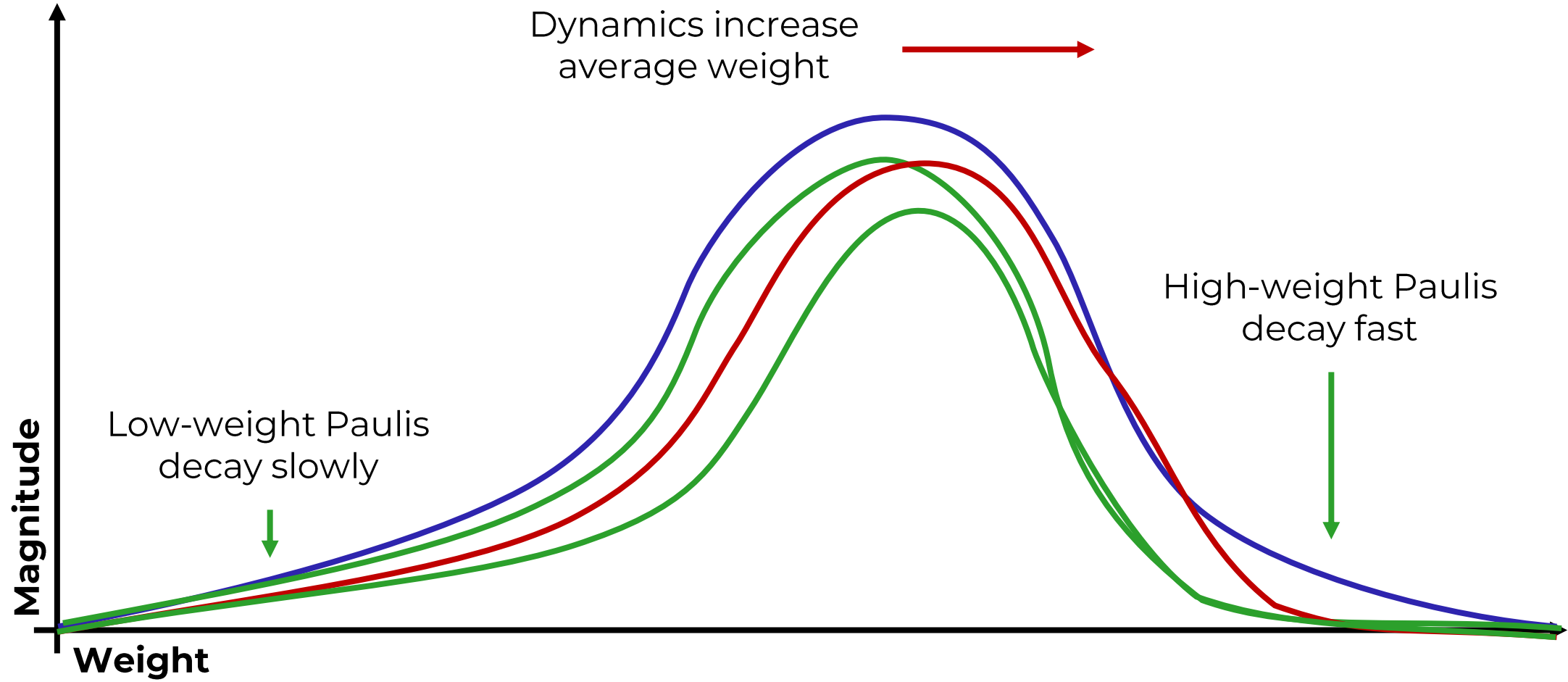
TOOL: PAULI EXPANSION

We can expand the system state
into Pauli words, there local
depolarizing noise acts like

$$\mathcal{D}_\lambda[P] = \lambda^{w(P)} P$$

¹Cleve et al., Quant. Inf. Comp. 16, 721–756 (2016)

Distribution of Pauli Coefficients



Application Takeaway

Error mitigation can still help,
but **error correction**
is what we really need

Technical Takeaway

We can prove much **stronger limitations** if we go **beyond generic arguments**

Part II

How does the picture change
with **finite samples**?

Quantum Metrology

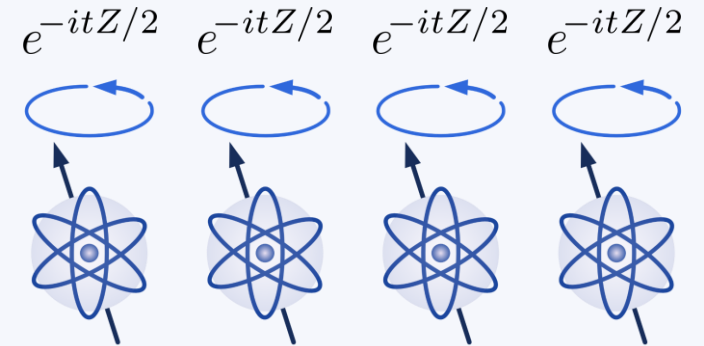
GOAL

Devise a protocol that estimates the phase as well as possible

PHASE ESTIMATION

Local evolution of an ensemble of spins under the phase Hamiltonian

$$H = \frac{1}{2} \sum_{i=1}^n Z_i$$



LIMITATIONS OF PRACTICAL METROLOGY

- › Low numbers of quantum systems
- › Slow operation speed of certain platforms
- › Drift of the system parameters



Practical metrology happens in the **finite-sample regime!**

Traditional Quantum Metrology

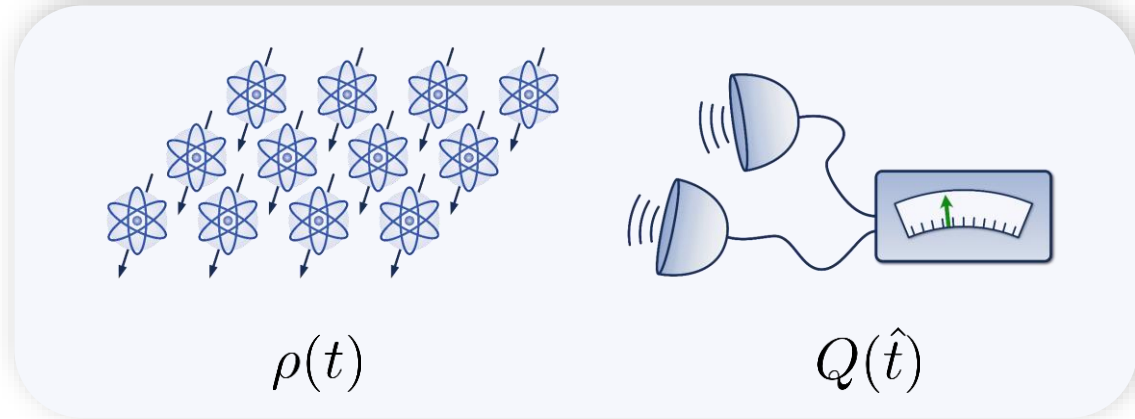


FIGURE OF MERIT

- › **Mean squared error** of the estimate
- › Equal to the **variance** if the estimate is exact in expectation

CRAMÉR-RAO BOUND¹

$$\text{Var}(\hat{t}) \geq \frac{1}{\mathcal{F}(t)}$$

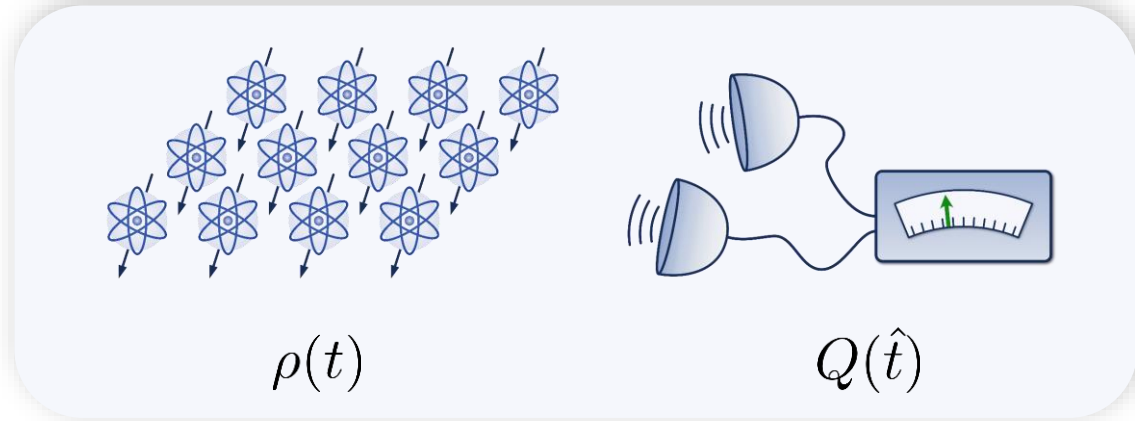
WHAT ABOUT THE FINITE-SAMPLE REGIME?

- › Cramér-Rao bound still constrains precision
- › But it can be **overly optimistic**

Optimizing for large Fisher information can lead to poor finite-sample performance

¹Helstrom, *Phys. Lett. A* (1967)

Single-shot Quantum Metrology



SINGLE-SHOT FIGURE OF MERIT

- › Define an **estimation tolerance** δ that should be achieved
- › Determine the probability of estimating the parameter within that tolerance

SUCCESS PROBABILITY¹

$$\mathbb{P}[|t - \hat{t}| \leq \delta]$$

¹See also Hayashi, *J. Phys. A* (2002), Walter and Renes, *IEEE Trans. Inform. Theory* (2014)

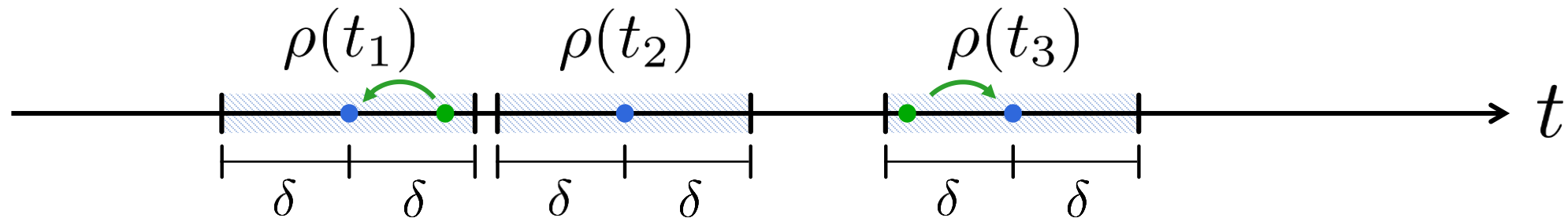
Hypothesis Testing Bound

MOTIVATION

If states that are $O(\delta)$ apart are hard to distinguish, estimating the parameter to precision δ should also be hard



Perform a reduction from quantum metrology to multi-hypothesis testing



We can use the metrology protocol to solve the multi-hypothesis testing task



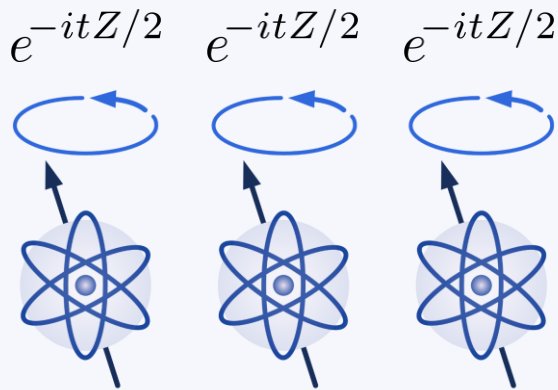
THEOREM

The success probability of quantum metrology cannot exceed the success probability of distinguishing states at times that are at least 2δ apart

Optimal Phase Estimation

PHASE ESTIMATION

Local evolution of an ensemble of spins under the phase Hamiltonian



$$H = \frac{1}{2} \sum_{i=1}^n Z_i$$

We show that the **pretty good measurement**¹ is optimal for covariant state sets



We use this to obtain a **closed-form solution** for the minimax success probability



The closed-form solution facilitates a **numerical comparison** of different **probe states**

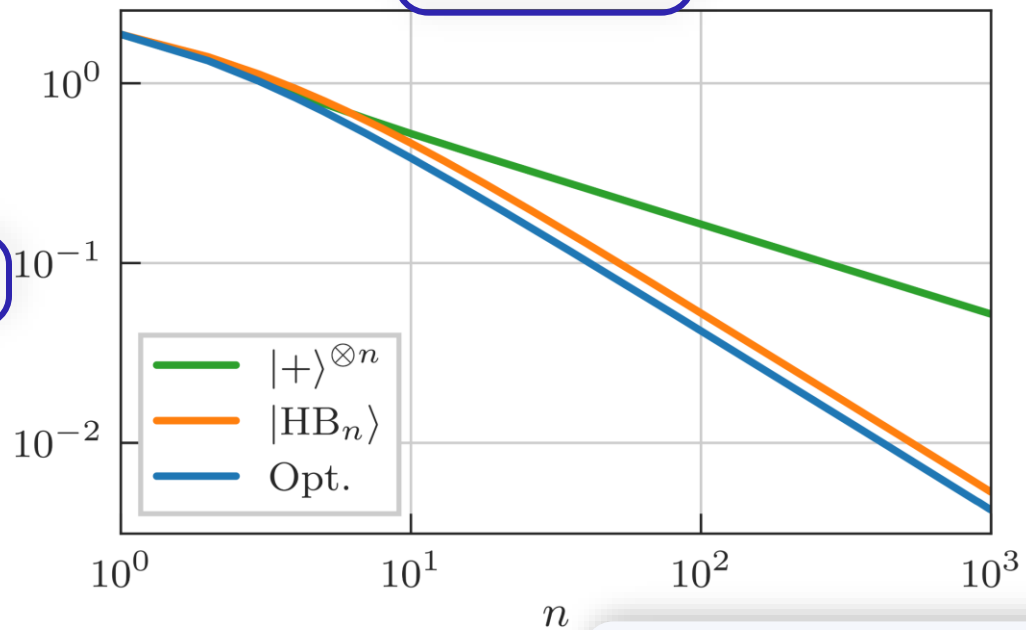
¹Holevo, *Rep. Math. Phys.* (1997)

Minimax Estimation Tolerance

Desired probability of success

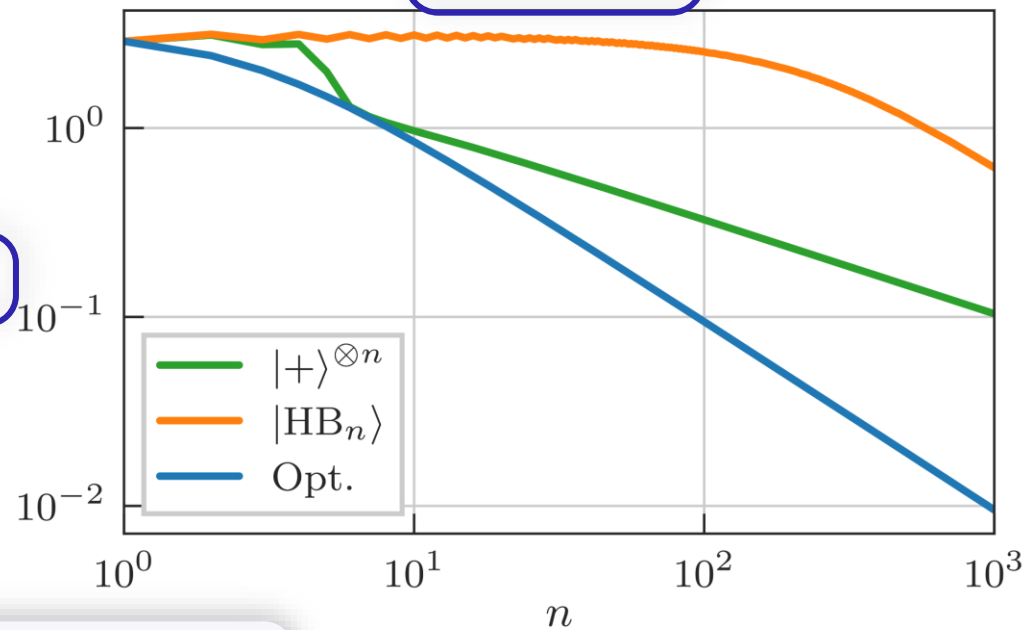
$\bar{\eta} = 0.9$

ϵ^*



$\bar{\eta} = 0.999$

ϵ^*



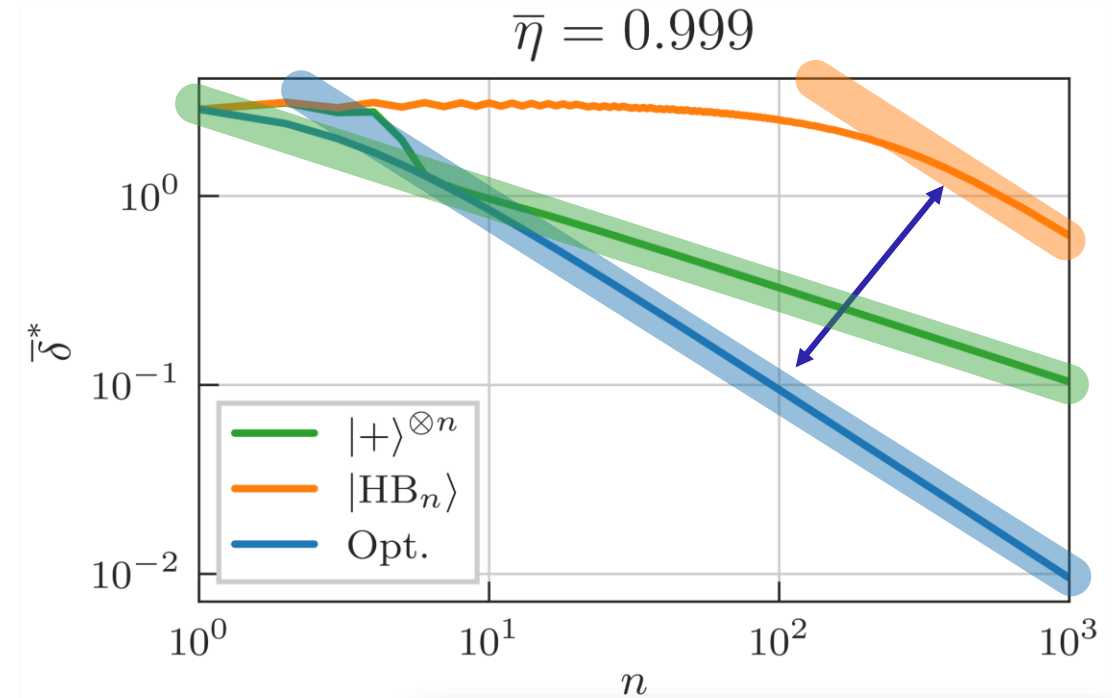
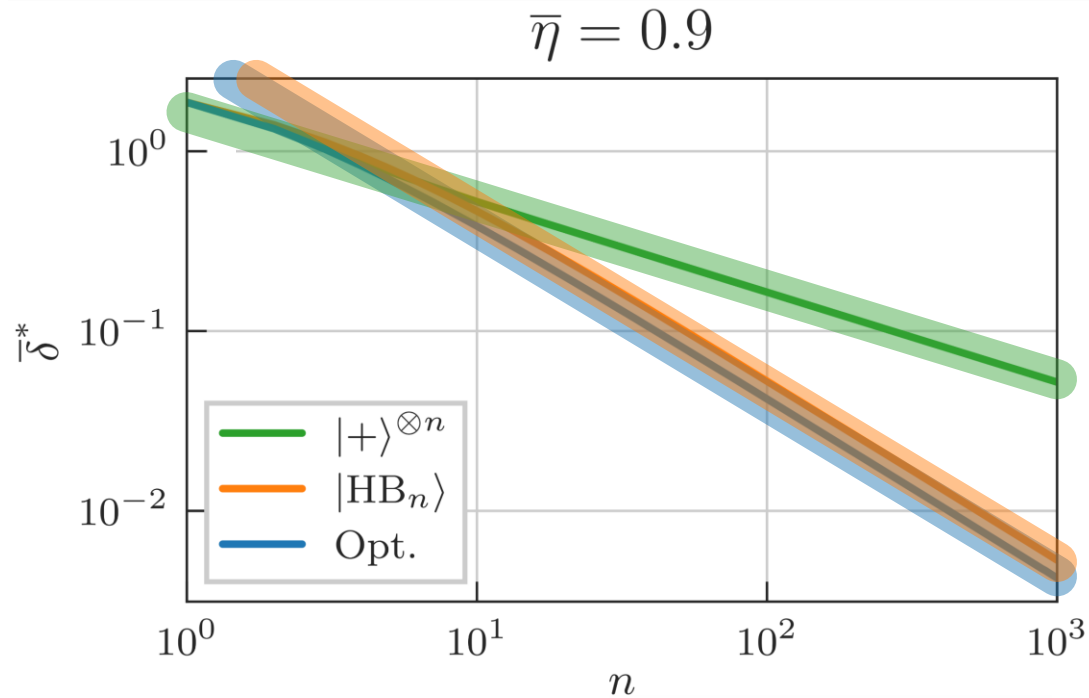
Achievable optimal tolerance
(Loglog plot)

Minimax Estimation Tolerance

Holland-Burnett State¹

$$|\text{HB}_n\rangle = \frac{1}{\sqrt{n+1}}(|0\rangle + |1\rangle + |2\rangle + \dots + |n\rangle)$$

Has similar Fisher information
as optimal state



Performance gap depending
on the success probability

**Cannot be predicted by Fisher
information!**

¹Holland and Burnett, *Phys. Rev. Lett.* (1993)

Beyond Quantum Fisher Information

The mean-squared error and the Cramér-Rao bound do not faithfully capture the finite-sample regime



Quantifying the **success probability** allows for a more fine-grained analysis



Allows us to rigorously show cases where the limits predicted by the Quantum Fisher Information can not be reached

Application Takeaway

Optimizing for Fisher
Information may lead to **poor
finite-sample performance**

Technical Takeaway

To fully understand limitations
we need to **consider the
choice of figure of merit**

Part III

How constraining are
computational restrictions?

Reviewing our Tools

Part I

Sample-complexity lower bounds based on Fano's inequality for estimator in hypothesis testing

Device Noise

$$\mathbb{P}[\hat{X} \neq X] \geq 1 - \frac{I(X : \hat{X}) + \log 2}{\log M}$$

Part II

Reduction to the error probability of a symmetric binary hypothesis test

Finite Samples

$$P_e^*(\rho \parallel \sigma) = \frac{1}{2} - \frac{1}{4} \|\rho - \sigma\|_1$$

Reviewing our Tools

Information-theoretic tools and bounds are indispensable when studying the fundamental limitations of quantum technologies



But these tools are often **asymptotic** in nature and assume **infinite computational power**



Can we come up with information-theoretic quantities that do not impose these assumptions?

Starting Point: Relative Entropy

RELATIVE ENTROPY

$$D(\rho \parallel \sigma) := \text{Tr}[\rho(\log \rho - \log \sigma)]$$

PARENT QUANTITY FOR

- › Entropy
- › Mutual Information
- › Coherent Information
- › Holevo Chi
- › ...

Single-shot Approach to Relative Entropy

HYPOTHESIS TESTING RELATIVE ENTROPY

Operational definition of the best type II error achievable at fixed type I error

$$D_h^\epsilon(\rho \parallel \sigma) := -\log \min_{0 \leq \Lambda \leq \mathbb{I}} \{ \text{Tr}[\Lambda\sigma] \mid \text{Tr}[\Lambda\rho] \geq 1 - \epsilon \}$$

QUANTUM STEIN'S LEMMA

Best asymptotic exponent in the limit of many copies

$$\lim_{n \rightarrow \infty} \frac{1}{n} D_h^\epsilon(\rho^{\otimes n} \parallel \sigma^{\otimes n}) = D(\rho \parallel \sigma)$$

Adding Complexity Restrictions

G-COMPLEXITY HYPOTHESIS TESTING RELATIVE ENTROPY¹

Operational definition of the best type II error achievable at fixed type I error and fixed gate complexity of the measurement

$$D_h^\epsilon(\rho \parallel \sigma; G) := -\log \min_{0 \leq \Lambda \leq \mathbb{I}} \{ \text{Tr}[\Lambda\sigma] \mid \text{Tr}[\Lambda\rho] \geq 1 - \epsilon, C(\Lambda) \leq G \}$$

COMPUTATIONAL RELATIVE ENTROPY

Regularizing this over polynomially many copies and gates gives the best exponent in the polynomial regime

$$\approx D_h^\epsilon(\rho^{\otimes \text{poly}(n)} \parallel \sigma^{\otimes \text{poly}(n)}; \text{poly}(n))$$

¹Munson et al., PRX Quantum 6 010346 (2025)

Teaser

- ▶ We define a computational relative entropy in a mathematically rigorous way
- ▶ We can show separations: there exists states that can be arbitrarily well tested without resource constraints, but not under computational constraints
- ▶ We derive a computational entropy and show that it has an interpretation as the best compression rate under computational constraints
- ▶ We show how our computational quantities can be used to understand entanglement manipulation under complexity constraints

Application Takeaway

Information-theoretic bounds
can be **arbitrarily far** from the
actual achievable
performance

Technical Takeaway

We can define **entropic quantities** that have **computational limitations** inbuilt

Conclusion

Lots of
exciting science
is left to be done

Exciting directions

- › We have a good understanding of information-theoretic limits on quantum technologies
- › But in reality, these applications are very restricted and can often not saturate these limits
- › Understanding the true limitations is thus still wide open, and we are still just at the beginning of fully capturing limitations that arise from
 - › Limited computational power
 - › Imperfect devices
 - › Limited sample sizes
 - › Limited capabilities (e.g. multi-copy measurements)

Thank you for your attention.

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Metrology



Mitigation



Slides